

Mean-field evolution of fermions with Coulomb interaction

Chiara Saffirio

Universität Zürich

in collaboration with
M. Porta, S. Rademacher and B. Schlein

Munich, March 30-April 1, 2017.

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L^2_a(\mathbb{R}^{3N})$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L^2_a(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L^2_a(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L^2_a(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

General goal: study the dynamics of the low energy states

$$\begin{cases} i\partial_t \psi_{N,t} = H_N \psi_{N,t}, \\ \psi_{N,0} = \psi_N. \end{cases}$$

N is large \implies look for **scaling regimes** in which the evolution can be well approximated by an **effective dynamics**.

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$



Scaling?

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

Space variable scaling: $X_j = N^{\frac{1}{3}} x_j$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

Space variable scaling: $X_j = N^{\frac{1}{3}} x_j$

$$\begin{aligned} H_N &= \sum_{j=1}^N \left[-N^{\frac{2}{3}} \Delta_{x_j} - \frac{N^{\frac{4}{3}}}{|x_j|} \right] + N^{\frac{1}{3}} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \\ &= N^{\frac{4}{3}} \left\{ \sum_{j=1}^N \left[-N^{-\frac{2}{3}} \Delta_{x_j} - \frac{1}{|x_j|} \right] + \frac{1}{N} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \right\} \end{aligned}$$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

Space variable scaling: $X_j = N^{\frac{1}{3}} x_j$

$$\begin{aligned} H_N &= \sum_{j=1}^N \left[-N^{\frac{2}{3}} \Delta_{x_j} - \frac{N^{\frac{4}{3}}}{|x_j|} \right] + N^{\frac{1}{3}} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \\ &= N^{\frac{4}{3}} \left\{ \sum_{j=1}^N \left[-\varepsilon^2 \Delta_{x_j} - \frac{1}{|x_j|} \right] + \frac{1}{N} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \right\} \end{aligned}$$

$$\varepsilon = N^{-\frac{1}{3}}$$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

Space variable scaling: $X_j = N^{\frac{1}{3}} x_j$

$$i\partial_t \psi_{N,t} = N^{\frac{4}{3}} \left\{ \sum_{j=1}^N \left[-\varepsilon^2 \Delta_{x_j} - \frac{1}{|x_j|} \right] + \frac{1}{N} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \right\} \psi_{N,t}$$

$$\varepsilon = N^{-\frac{1}{3}}$$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

Space variable scaling: $X_j = N^{\frac{1}{3}} x_j$

$$i\partial_t \psi_{N,t} = N^{\frac{4}{3}} \left\{ \sum_{j=1}^N \left[-\varepsilon^2 \Delta_{x_j} - \frac{1}{|x_j|} \right] + \frac{1}{N} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \right\} \psi_{N,t}$$

Time scaling:

$$i\varepsilon \partial_t \psi_{N,t} = \left\{ \sum_{j=1}^N \left[-\varepsilon^2 \Delta_{x_j} - \frac{1}{|x_j|} \right] + \frac{1}{N} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \right\} \psi_{N,t} \quad \varepsilon = N^{-\frac{1}{3}}$$

Setting

Consider a system of N interacting fermions with wave function $\psi_N \in L_a^2(\mathbb{R}^{3N})$

Interaction: $V(x) = \frac{1}{|x|}$

Motivation: dynamics of large atoms (e.g. electrically neutral atom)

$$H_N = \sum_{j=1}^N \left[-\Delta_{x_j} - \frac{N}{|x_j|} \right] + \sum_{i < j}^N \frac{1}{|x_i - x_j|}$$

Space variable scaling: $X_j = N^{\frac{1}{3}} x_j$

$$i\partial_t \psi_{N,t} = N^{\frac{4}{3}} \left\{ \sum_{j=1}^N \left[-\varepsilon^2 \Delta_{x_j} - \frac{1}{|x_j|} \right] + \frac{1}{N} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \right\} \psi_{N,t}$$

Time scaling:

$$i\varepsilon \partial_t \psi_{N,t} = \left\{ \sum_{j=1}^N \left[-\varepsilon^2 \Delta_{x_j} - \frac{1}{|x_j|} \right] + \frac{1}{N} \sum_{i < j}^N \frac{1}{|x_i - x_j|} \right\} \psi_{N,t} \quad \varepsilon = N^{-\frac{1}{3}}$$

Effective evolution equation?

Ground state \simeq Slater determinant (Bach '92, Graf & Solovej '94)

$$\psi_{\text{Slater}}(\mathbf{x}_1, \dots, \mathbf{x}_N) = \frac{1}{\sqrt{N!}} \det(f_j(\mathbf{x}_i))_{i,j \leq N}, \quad \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

Effective evolution equation?

Ground state \simeq Slater determinant (Bach '92, Graf & Solovej '94)

$$\psi_{\text{Slater}}(\mathbf{x}_1, \dots, \mathbf{x}_N) = \frac{1}{\sqrt{N!}} \det(f_j(\mathbf{x}_i))_{i,j \leq N}, \quad \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

Reduced one-particle density matrix:

$$\omega_N = N \operatorname{tr}_{2 \dots N} |\psi_{\text{Slater}}\rangle \langle \psi_{\text{Slater}}| = \sum_{j=1}^N |f_j\rangle \langle f_j|$$

Effective evolution equation?

Ground state \simeq Slater determinant (Bach '92, Graf & Solovej '94)

$$\psi_{\text{Slater}}(\mathbf{x}_1, \dots, \mathbf{x}_N) = \frac{1}{\sqrt{N!}} \det(f_j(\mathbf{x}_i))_{i,j \leq N}, \quad \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

Reduced one-particle density matrix:

$$\omega_N = N \operatorname{tr}_{2 \dots N} |\psi_{\text{Slater}}\rangle \langle \psi_{\text{Slater}}| = \sum_{j=1}^N |f_j\rangle \langle f_j| \text{ minimiser of the HF energy}$$

$$\mathcal{E}_{\text{HF}}(\omega_N) = \operatorname{Tr}(-\Delta + V_{\text{ext}})\omega_N + \frac{1}{2N^{1/3}} \int dx dy \frac{1}{|x-y|} [\rho(x)\rho(y) - |\omega_N(x;y)|^2]$$

$$\text{where } \rho(x) = \frac{1}{N} \omega_N(x;x).$$

Effective evolution equation?

Ground state \simeq Slater determinant (Bach '92, Graf & Solovej '94)

$$\psi_{\text{Slater}}(\mathbf{x}_1, \dots, \mathbf{x}_N) = \frac{1}{\sqrt{N!}} \det(f_j(x_i))_{i,j \leq N}, \quad \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

Reduced one-particle density matrix:

$$\omega_N = N \operatorname{tr}_{2 \dots N} |\psi_{\text{Slater}}\rangle \langle \psi_{\text{Slater}}| = \sum_{j=1}^N |f_j\rangle \langle f_j| \text{ minimiser of the HF energy}$$

$$\mathcal{E}_{\text{HF}}(\omega_N) = \operatorname{Tr}(-\Delta + V_{\text{ext}})\omega_N + \frac{1}{2N^{1/3}} \int dx dy \frac{1}{|x-y|} [\rho(x)\rho(y) - |\omega_N(x;y)|^2]$$

where $\rho(x) = \frac{1}{N} \omega_N(x;x)$.

$$\text{Dynamics} \implies \gamma_{N,t}^{(1)} = N \operatorname{tr}_{2 \dots N} |\psi_{N,t}\rangle \langle \psi_{N,t}| \simeq \omega_{N,t}$$

Effective evolution equation?

Ground state \simeq Slater determinant (Bach '92, Graf & Solovej '94)

$$\psi_{\text{Slater}}(\mathbf{x}_1, \dots, \mathbf{x}_N) = \frac{1}{\sqrt{N!}} \det(f_j(x_i))_{i,j \leq N}, \quad \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

Reduced one-particle density matrix:

$$\omega_N = N \operatorname{tr}_{2 \dots N} |\psi_{\text{Slater}}\rangle \langle \psi_{\text{Slater}}| = \sum_{j=1}^N |f_j\rangle \langle f_j| \text{ minimiser of the HF energy}$$

$$\mathcal{E}_{\text{HF}}(\omega_N) = \operatorname{Tr}(-\Delta + V_{\text{ext}})\omega_N + \frac{1}{2N^{1/3}} \int dx dy \frac{1}{|x-y|} [\rho(x)\rho(y) - |\omega_N(x; y)|^2]$$

where $\rho(x) = \frac{1}{N} \omega_N(x; x)$.

Dynamics $\implies \gamma_{N,t}^{(1)} = N \operatorname{tr}_{2 \dots N} |\psi_{N,t}\rangle \langle \psi_{N,t}| \simeq \omega_{N,t}$

$\omega_{N,t}$ solution to the time dependent **Hartree-Fock eqn**

$$i\varepsilon \partial_t \omega_{N,t} = [-\varepsilon^2 \Delta_{x_j} + \frac{1}{|\cdot|} * \rho_t - X_t, \omega_{N,t}]$$

with $X_t(x; y) = \frac{1}{N} \frac{1}{|x-y|} \omega_{N,t}(x; y)$.

Effective evolution equation?

Ground state \simeq Slater determinant (Bach '92, Graf & Solovej '94)

$$\psi_{\text{Slater}}(\mathbf{x}_1, \dots, \mathbf{x}_N) = \frac{1}{\sqrt{N!}} \det(f_j(x_i))_{i,j \leq N}, \quad \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

Reduced one-particle density matrix:

$$\omega_N = N \operatorname{tr}_{2 \dots N} |\psi_{\text{Slater}}\rangle \langle \psi_{\text{Slater}}| = \sum_{j=1}^N |f_j\rangle \langle f_j| \text{ minimiser of the HF energy}$$

$$\mathcal{E}_{\text{HF}}(\omega_N) = \operatorname{Tr}(-\Delta + V_{\text{ext}})\omega_N + \frac{1}{2N^{1/3}} \int dx dy \frac{1}{|x-y|} [\rho(x)\rho(y) - |\omega_N(x; y)|^2]$$

$$\text{where } \rho(x) = \frac{1}{N} \omega_N(x; x).$$

$$\text{Dynamics} \implies \gamma_{N,t}^{(1)} = N \operatorname{tr}_{2 \dots N} |\psi_{N,t}\rangle \langle \psi_{N,t}| \simeq \omega_{N,t}$$

$\omega_{N,t}$ solution to the time dependent **Hartree-Fock eqn**

$$i\varepsilon \partial_t \omega_{N,t} = [-\varepsilon^2 \Delta_{x_j} + \frac{1}{|\cdot|} * \rho_t - X_t, \omega_{N,t}]$$

$$\text{with } X_t(x; y) = \frac{1}{N} \frac{1}{|x-y|} \omega_{N,t}(x; y).$$

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Assume there exist $T > 0$, $p > 5$, $C > 0$ such that $\forall t \in [0, T]$

$$\|\rho|_{[x, \omega_{N,t}]}\|_{L^1} \leq CN\varepsilon, \quad \|\rho|_{[x, \omega_{N,t}]}\|_{L^p} \leq CN\varepsilon.$$

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Assume there exist $T > 0$, $p > 5$, $C > 0$ such that $\forall t \in [0, T]$

$$\|\rho|_{[X, \omega_{N,t}]}\|_{L^1} \leq CN\varepsilon, \quad \|\rho|_{[X, \omega_{N,t}]}\|_{L^p} \leq CN\varepsilon.$$

Let $\psi_N \in L^2_{\text{a}}(\mathbb{R}^{3N})$ be s.t. its one-particle reduced density matrix $\gamma_N^{(1)}$ satisfies

$$\text{tr } |\gamma_N^{(1)} - \omega_N| \leq CN^\alpha, \quad \alpha \in [0, 1).$$

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Assume there exist $T > 0$, $p > 5$, $C > 0$ such that $\forall t \in [0, T]$

$$\|\rho|_{[x, \omega_{N,t}]}\|_{L^1} \leq CN\varepsilon, \quad \|\rho|_{[x, \omega_{N,t}]}\|_{L^p} \leq CN\varepsilon.$$

Let $\psi_N \in L^2_{\text{a}}(\mathbb{R}^{3N})$ be s.t. its one-particle reduced density matrix $\gamma_N^{(1)}$ satisfies

$$\text{tr } |\gamma_N^{(1)} - \omega_N| \leq CN^\alpha, \quad \alpha \in [0, 1).$$

Consider the evolution $\psi_{N,t} = e^{-itH_N/\varepsilon} \psi_N$ and $\gamma_{N,t}^{(1)}$ the corresponding one-particle reduced density matrix.

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Assume there exist $T > 0$, $p > 5$, $C > 0$ such that $\forall t \in [0, T]$

$$\|\rho|_{[x, \omega_{N,t}]}\|_{L^1} \leq CN\varepsilon, \quad \|\rho|_{[x, \omega_{N,t}]}\|_{L^p} \leq CN\varepsilon.$$

Let $\psi_N \in L^2_a(\mathbb{R}^{3N})$ be s.t. its one-particle reduced density matrix $\gamma_N^{(1)}$ satisfies

$$\text{tr } |\gamma_N^{(1)} - \omega_N| \leq CN^\alpha, \quad \alpha \in [0, 1).$$

Consider the evolution $\psi_{N,t} = e^{-itH_N/\varepsilon} \psi_N$ and $\gamma_{N,t}^{(1)}$ the corresponding one-particle reduced density matrix.

Then for every $\delta > 0$, there exists $C > 0$ s.t.

$$\text{tr } |\gamma_{N,t}^{(1)} - \omega_{N,t}| \leq C(t)(N^\alpha + N^{11/12+\delta}).$$

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Assume there exist $T > 0$, $p > 5$, $C > 0$ such that $\forall t \in [0, T]$

$$\|\rho|_{[x, \omega_{N,t}]}\|_{L^1} \leq CN\varepsilon, \quad \|\rho|_{[x, \omega_{N,t}]}\|_{L^p} \leq CN\varepsilon.$$

Let $\psi_N \in L^2_a(\mathbb{R}^{3N})$ be s.t. its one-particle reduced density matrix $\gamma_N^{(1)}$ satisfies

$$\text{tr } |\gamma_N^{(1)} - \omega_N| \leq CN^\alpha, \quad \alpha \in [0, 1).$$

Consider the evolution $\psi_{N,t} = e^{-itH_N/\varepsilon} \psi_N$ and $\gamma_{N,t}^{(1)}$ the corresponding one-particle reduced density matrix.

Then for every $\delta > 0$, there exists $C > 0$ s.t.

$$\text{tr } |\gamma_{N,t}^{(1)} - \omega_{N,t}| \leq C(t)(N^\alpha + N^{11/12+\delta}).$$

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Assume there exist $T > 0$, $p > 5$, $C > 0$ such that $\forall t \in [0, T]$

$$\|\rho_{|[x, \omega_{N,t}]}\|_{L^1} \leq CN\varepsilon, \quad \|\rho_{|[x, \omega_{N,t}]}\|_{L^p} \leq CN\varepsilon.$$

Let $\psi_N \in L^2_a(\mathbb{R}^{3N})$ be s.t. its one-particle reduced density matrix $\gamma_N^{(1)}$ satisfies

$$\text{tr } |\gamma_N^{(1)} - \omega_N| \leq CN^\alpha, \quad \alpha \in [0, 1).$$

Consider the evolution $\psi_{N,t} = e^{-itH_N/\varepsilon} \psi_N$ and $\gamma_{N,t}^{(1)}$ the corresponding one-particle reduced density matrix.

Then for every $\delta > 0$, there exists $C > 0$ s.t.

$$\text{tr } |\gamma_{N,t}^{(1)} - \omega_{N,t}| \leq C(t)(N^\alpha + N^{11/12+\delta}).$$

Our result

Theorem

Let ω_N be a sequence of orthogonal projections on $L^2(\mathbb{R}^3)$ with $\text{tr } \omega_N = N$ and $\text{tr } (-\varepsilon^2 \Delta) \omega_N \leq CN$.

Let $\omega_{N,t}$ be a solution to the TDHF equation with initial data ω_N .

Assume there exist $T > 0$, $p > 5$, $C > 0$ such that $\forall t \in [0, T]$

$$\|\rho|_{[x, \omega_{N,t}]}\|_{L^1} \leq CN\varepsilon, \quad \|\rho|_{[x, \omega_{N,t}]}\|_{L^p} \leq CN\varepsilon.$$

Let $\psi_N \in L^2_a(\mathbb{R}^{3N})$ be s.t. its one-particle reduced density matrix $\gamma_N^{(1)}$ satisfies

$$\text{tr } |\gamma_N^{(1)} - \omega_N| \leq CN^\alpha, \quad \alpha \in [0, 1).$$

Consider the evolution $\psi_{N,t} = e^{-itH_N/\varepsilon} \psi_N$ and $\gamma_{N,t}^{(1)}$ the corresponding one-particle reduced density matrix.

Then for every $\delta > 0$, there exists $C > 0$ s.t.

$$\text{tr } |\gamma_{N,t}^{(1)} - \omega_{N,t}| \leq C(t)(N^\alpha + N^{11/12+\delta}).$$

Remarks

- Reduce the many-body problem to a PDE problem

Remarks

- Reduce the many-body problem to a PDE problem
- Translation invariant states

Remarks

- Reduce the many-body problem to a PDE problem
- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

Remarks

- Reduce the many-body problem to a PDE problem

- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

$[X, \omega_{N,t}]$ and $|[X, \omega_{N,t}]|$ translation invariant

Remarks

- Reduce the many-body problem to a PDE problem
- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

$[X, \omega_{N,t}]$ and $|[X, \omega_{N,t}]|$ translation invariant $\implies \rho_{|[X, \omega_{N,t}]|} \equiv \text{const}$

Remarks

- Reduce the many-body problem to a PDE problem
- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

$[x, \omega_{N,t}]$ and $|[x, \omega_{N,t}]|$ translation invariant $\implies \rho_{|[x, \omega_{N,t}]|} \equiv \text{const}$

$$\rho_{|[x, \omega_N]|} = O(N\varepsilon)$$

ω_N function of $(x - y)$ decaying at distance $|x - y| \gg \varepsilon$.

Remarks

- Reduce the many-body problem to a PDE problem

- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

$[\chi, \omega_{N,t}]$ and $||[\chi, \omega_{N,t}]||$ translation invariant $\implies \rho_{|[\chi, \omega_{N,t}]|} \equiv \text{const}$

$$\rho_{|[\chi, \omega_N]|} = O(N\varepsilon)$$

ω_N function of $(x - y)$ decaying at distance $|x - y| \gg \varepsilon$.

- N dependence in the HF equation

Remarks

- Reduce the many-body problem to a PDE problem

- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

$[\chi, \omega_{N,t}]$ and $||[\chi, \omega_{N,t}]||$ translation invariant $\implies \rho_{|[\chi, \omega_{N,t}]|} \equiv \text{const}$

$$\rho_{|[\chi, \omega_N]|} = O(N\varepsilon)$$

ω_N function of $(x - y)$ decaying at distance $|x - y| \gg \varepsilon$.

- N dependence in the HF equation: $N \rightarrow \infty$?

Remarks

- Reduce the many-body problem to a PDE problem

- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

$[x, \omega_{N,t}]$ and $|[x, \omega_{N,t}]|$ translation invariant $\implies \rho_{|[x, \omega_{N,t}]|} \equiv \text{const}$

$$\rho_{|[x, \omega_N]|} = O(N\varepsilon)$$

ω_N function of $(x - y)$ decaying at distance $|x - y| \gg \varepsilon$.

- N dependence in the HF equation: $N \rightarrow \infty$?
The **Vlasov equation** is the next degree of approximation.

Remarks

- Reduce the many-body problem to a PDE problem

- Translation invariant states

Λ box with periodic b.c. $\implies \omega_{N,t}(x; y) \simeq \omega_{N,t}(x - y)$

$[x, \omega_{N,t}]$ and $|[x, \omega_{N,t}]|$ translation invariant $\implies \rho_{|[x, \omega_{N,t}]|} \equiv \text{const}$

$$\rho_{|[x, \omega_N]|} = O(N\varepsilon)$$

ω_N function of $(x - y)$ decaying at distance $|x - y| \gg \varepsilon$.

- N dependence in the HF equation: $N \rightarrow \infty$?

The **Vlasov equation** is the next degree of approximation.

The Wigner transform of $\omega_{N,t}$ for $N \rightarrow \infty$ solves the Vlasov equation

$$\partial_t W_t(x, v) + v \cdot \nabla_x W_t(x, v) = \nabla(V * \rho_t)(x) \cdot \nabla_v W_t(x, v)$$

Previous results

- 1980: Narnhofer and Sewell (convergence to Vlasov for analytic potentials);
- 1982: Spohn (extension to C^2 potentials);

Previous results

- 1980: Narnhofer and Sewell (convergence to Vlasov for analytic potentials);
- 1982: Spohn (extension to C^2 potentials);
- 2003: Elgart, Erdős, Schlein, Yau (TDHF short times, analytic pot.);

Previous results

- 1980: Narnhofer and Sewell (convergence to Vlasov for analytic potentials);
- 1982: Spohn (extension to C^2 potentials);
- 2003: Elgart, Erdős, Schlein, Yau (TDHF short times, analytic pot.);
- other scalings:
 - ▶ 2002: Bardos, Golse, Gottlieb, Mauser;
 - ▶ 2010: Fröhlich, Knowles;

Previous results

- 1980: Narnhofer and Sewell (convergence to Vlasov for analytic potentials);
- 1982: Spohn (extension to C^2 potentials);
- 2003: Elgart, Erdős, Schlein, Yau (TDHF short times, analytic pot.);
- other scalings:
 - ▶ 2002: Bardos, Golse, Gottlieb, Mauser;
 - ▶ 2010: Fröhlich, Knowles;
- 2013: Benedikter, Porta, Schlein (larger class of pot., all times);
- 2015: Benedikter, Jaksic, Porta, S., Schlein (mixed states);

Previous results

- 1980: Narnhofer and Sewell (convergence to Vlasov for analytic potentials);
- 1982: Spohn (extension to C^2 potentials);
- 2003: Elgart, Erdős, Schlein, Yau (TDHF short times, analytic pot.);
- other scalings:
 - ▶ 2002: Bardos, Golse, Gottlieb, Mauser;
 - ▶ 2010: Fröhlich, Knowles;
- 2013: Benedikter, Porta, Schlein (larger class of pot., all times);
- 2015: Benedikter, Jaksic, Porta, S., Schlein (mixed states);
- 2016 Coulomb potential with different scalings:
 - ▶ Bach, Breteaux, Petrat, Pickl, Tzaneteas;
 - ▶ Petrat, Pickl.

Idea of the proof

- **Fock space representation:**

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

Idea of the proof

- **Fock space representation:**

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

creation and *annihilation* operators:

Idea of the proof

- **Fock space representation:**

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

creation and *annihilation* operators: $\forall f \in L^2(\mathbb{R}^3)$

Idea of the proof

- Fock space representation:

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

creation and *annihilation* operators: $\forall f \in L^2(\mathbb{R}^3)$

$$a^*(f) = \int dx f(x) a_x^* \quad a(f) = \int dx \overline{f(x)} a_x$$

Idea of the proof

- Fock space representation:

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

creation and *annihilation* operators: $\forall f \in L^2(\mathbb{R}^3)$

$$a^*(f) = \int dx f(x) a_x^* \quad a(f) = \int dx \overline{f(x)} a_x$$

number of particle operator:

$$\mathcal{N} = \int dx a_x^* a_x$$

Idea of the proof

- **Fock space representation:**

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

creation and *annihilation* operators: $\forall f \in L^2(\mathbb{R}^3)$

$$a^*(f) = \int dx f(x) a_x^* \quad a(f) = \int dx \overline{f(x)} a_x$$

number of particle operator:

$$\mathcal{N} = \int dx a_x^* a_x$$

second quantisation of J , operator on $L^2(\mathbb{R}^3)$

$$d\Gamma(J) = \iint dx dy J(x; y) a_x^* a_y$$

Idea of the proof

- **Fock space representation:**

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

creation and *annihilation* operators: $\forall f \in L^2(\mathbb{R}^3)$

$$a^*(f) = \int dx f(x) a_x^* \quad a(f) = \int dx \overline{f(x)} a_x$$

number of particle operator:

$$\mathcal{N} = \int dx a_x^* a_x$$

second quantisation of J , operator on $L^2(\mathbb{R}^3)$

$$d\Gamma(J) = \iint dx dy J(x; y) a_x^* a_y$$

Let $\Psi \in \mathcal{F}$,

$$\gamma_\Psi^{(1)}(x; y) = \langle \Psi, a_x^* a_y \Psi \rangle$$

Idea of the proof

- **Fock space representation:**

$$\mathcal{F} = \bigoplus_{n \geq 0} L_a^2(\mathbb{R}^{3n})$$

creation and *annihilation* operators: $\forall f \in L^2(\mathbb{R}^3)$

$$a^*(f) = \int dx f(x) a_x^* \quad a(f) = \int dx \overline{f(x)} a_x$$

number of particle operator:

$$\mathcal{N} = \int dx a_x^* a_x$$

second quantisation of J , operator on $L^2(\mathbb{R}^3)$

$$d\Gamma(J) = \iint dx dy J(x; y) a_x^* a_y$$

Let $\Psi \in \mathcal{F}$,

$$\gamma_\Psi^{(1)}(x; y) = \langle \Psi, a_x^* a_y \Psi \rangle \implies \text{tr } \gamma_\Psi^{(1)} = \langle \Psi, \mathcal{N} \Psi \rangle$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

$$\text{s.t. } \forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

s.t. $\forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$

- Compute $\gamma_{N,t}^{(1)}$:

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

$$\text{s.t. } \forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$$

- Compute $\gamma_{N,t}^{(1)}$:

$$\gamma_{N,t}^{(1)}(x; y) = \langle e^{-i\mathcal{H}_{Nt}/\varepsilon} R_{\omega_N} \Omega, a_x^* a_y e^{-i\mathcal{H}_{Nt}/\varepsilon} R_{\omega_N} \Omega \rangle$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

$$\text{s.t. } \forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$$

- Compute $\gamma_{N,t}^{(1)}$:

$$\gamma_{N,t}^{(1)}(x; y) = \langle R_{\omega_{N,t}} R_{\omega_{N,t}}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega, a_x^* a_y R_{\omega_{N,t}} R_{\omega_{N,t}}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega \rangle$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

$$\text{s.t. } \forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$$

- Compute $\gamma_{N,t}^{(1)}$:

$$\begin{aligned} \gamma_{N,t}^{(1)}(x; y) &= \langle R_{\omega_N,t} R_{\omega_N,t}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega, a_x^* a_y R_{\omega_N,t} R_{\omega_N,t}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, R_{\omega_N,t}^* a_x^* a_y R_{\omega_N,t} \mathcal{U}_N(t) \Omega \rangle \end{aligned}$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

s.t. $\forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$

- Compute $\gamma_{N,t}^{(1)}$:

$$\begin{aligned} \gamma_{N,t}^{(1)}(x; y) &= \langle R_{\omega_N,t} R_{\omega_N,t}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega, a_x^* a_y R_{\omega_N,t} R_{\omega_N,t}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, R_{\omega_N,t}^* a_x^* R_{\omega_N,t} a_y R_{\omega_N,t} \mathcal{U}_N(t) \Omega \rangle \end{aligned}$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

s.t. $\forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g)$.

- Compute $\gamma_{N,t}^{(1)}$:

$$\begin{aligned} \gamma_{N,t}^{(1)}(x; y) &= \langle R_{\omega_{N,t}} R_{\omega_{N,t}}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega, a_x^* a_y R_{\omega_{N,t}} R_{\omega_{N,t}}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, R_{\omega_{N,t}}^* a_x^* R_{\omega_{N,t}} R_{\omega_{N,t}}^* a_y R_{\omega_{N,t}} \mathcal{U}_N(t) \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, [a^*(1 - \omega_{t,x}) + a(\omega_{t,x})][a(1 - \omega_{t,y}) + a^*(\omega_{t,y})] \mathcal{U}_N(t) \Omega \rangle \end{aligned}$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

s.t. $\forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$

- Compute $\gamma_{N,t}^{(1)}$:

$$\begin{aligned} \gamma_{N,t}^{(1)}(x; y) &= \langle R_{\omega_{N,t}} R_{\omega_{N,t}}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega, a_x^* a_y R_{\omega_{N,t}} R_{\omega_{N,t}}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, R_{\omega_{N,t}}^* a_x^* R_{\omega_{N,t}} R_{\omega_{N,t}}^* a_y R_{\omega_{N,t}} \mathcal{U}_N(t) \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, [a^*(1 - \omega_{t,x}) + a(\omega_{t,x})][a(1 - \omega_{t,y}) + a^*(\omega_{t,y})] \mathcal{U}_N(t) \Omega \rangle \\ &= \omega_{N,t}(x; y) + \text{terms in normal order} \end{aligned}$$

Idea of the proof

- **Bogoliubov transformation:**

$$\omega_N = \sum_{j=1}^N |f_j\rangle\langle f_j|, \{f_j\}_{j \leq N} \text{ o.n.s. in } L^2(\mathbb{R}^3).$$

$R_{\omega_N} : \mathcal{F} \rightarrow \mathcal{F}$ unitary implementor of a Bogoliubov transformation

$$R_{\omega_N} \Omega = a^*(f_1) \dots a^*(f_N) \Omega$$

$$\text{s.t. } \forall g \in L^2(\mathbb{R}^3), R_{\omega_N}^* a^*(g) R_{\omega_N} = a^*((1 - \omega_N)g) + a(\omega_N g).$$

- Compute $\gamma_{N,t}^{(1)}$:

$$\begin{aligned} \gamma_{N,t}^{(1)}(x; y) &= \langle R_{\omega_N,t} R_{\omega_N,t}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega, a_x^* a_y R_{\omega_N,t} R_{\omega_N,t}^* e^{-i\mathcal{H}_N t/\varepsilon} R_{\omega_N} \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, R_{\omega_N,t}^* a_x^* R_{\omega_N,t} R_{\omega_N,t}^* a_y R_{\omega_N,t} \mathcal{U}_N(t) \Omega \rangle \\ &= \langle \mathcal{U}_N(t) \Omega, [a^*(1 - \omega_{t,x}) + a(\omega_{t,x})][a(1 - \omega_{t,y}) + a^*(\omega_{t,y})] \mathcal{U}_N(t) \Omega \rangle \\ &= \omega_{N,t}(x; y) + \text{terms in normal order} \end{aligned}$$

$$\text{tr} |\gamma_{N,t}^{(1)} - \omega_{N,t}| \lesssim \langle \mathcal{U}_N(t) \Omega, \mathcal{N} \mathcal{U}_N(t) \Omega \rangle$$

Idea of the proof

- **Control on the fluctuations**

Bound on the expected number of excitations of the Slater determinant

$$\langle \mathcal{U}_N(t)\Omega, \mathcal{N}\mathcal{U}_N(t)\Omega \rangle$$

Idea of the proof

- **Control on the fluctuations**

Bound on the expected number of excitations of the Slater determinant

$$\langle \mathcal{U}_N(t)\Omega, \mathcal{N}\mathcal{U}_N(t)\Omega \rangle$$

To control fluctuation: Grönwall type estimate

$$i\varepsilon \frac{d}{dt} \langle \mathcal{U}_N(t)\Omega, \mathcal{N}\mathcal{U}_N(t)\Omega \rangle = \dots$$

Idea of the proof

- **Control on the fluctuations**

Bound on the expected number of excitations of the Slater determinant

$$\langle \mathcal{U}_N(t)\Omega, \mathcal{N}\mathcal{U}_N(t)\Omega \rangle$$

To control fluctuation: Grönwall type estimate

$$\begin{aligned} & i\varepsilon \frac{d}{dt} \langle \mathcal{U}_N(t)\Omega, \mathcal{N}\mathcal{U}_N(t)\Omega \rangle \\ &= 4i \Im \frac{1}{N} \iint dx dy \frac{1}{|x-y|} \\ & \{ \langle \mathcal{U}_N(t)\Omega, \mathbf{a}(\omega_{t,x})\mathbf{a}(\omega_{t,y})\mathbf{a}(1-\omega_{t,y})\mathbf{a}(1-\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle \\ & + \langle \mathcal{U}_N(t)\Omega, \mathbf{a}^*(1-\omega_{t,y})\mathbf{a}^*(\omega_{t,y})\mathbf{a}^*(\omega_{t,x})\mathbf{a}(\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle \\ & + \langle \mathcal{U}_N(t)\Omega, \mathbf{a}^*(1-\omega_{t,x})\mathbf{a}(\omega_{t,y})\mathbf{a}(1-\omega_{t,y})\mathbf{a}(1-\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle \} \end{aligned}$$

Fefferman – de la Llave representation

$$(\star) \quad \frac{1}{N} \int dx dy \frac{1}{|x-y|} \langle \mathcal{U}_N(t)\Omega, a(\omega_{t,x})a(\omega_{t,y})a(1-\omega_{t,y})a(1-\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle$$

Fefferman – de la Llave representation

$$(\star) \quad \frac{1}{N} \int dx dy \frac{1}{|x-y|} \langle \mathcal{U}_N(t) \Omega, a(\omega_{t,x}) a(\omega_{t,y}) a(1-\omega_{t,y}) a(1-\omega_{t,x}) \mathcal{U}_N(t) \Omega \rangle$$

Smooth version of **Fefferman–de la Llave representation** for the Coulomb potential:

$$\frac{1}{|x-y|} = C \int_0^\infty \frac{dr}{r^5} \int dz \chi_{(r,z)}(x) \chi_{(r,z)}(y), \quad \chi_{(r,z)}(x) = e^{-|x-z|^2/r^2}$$

Fefferman – de la Llave representation

$$(*) \quad \frac{1}{N} \int dx dy \frac{1}{|x-y|} \langle \mathcal{U}_N(t)\Omega, a(\omega_{t,x})a(\omega_{t,y})a(1-\omega_{t,y})a(1-\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle$$

Smooth version of **Fefferman–de la Llave representation** for the Coulomb potential:

$$\frac{1}{|x-y|} = C \int_0^\infty \frac{dr}{r^5} \int dz \chi_{(r,z)}(x) \chi_{(r,z)}(y), \quad \chi_{(r,z)}(x) = e^{-|x-z|^2/r^2}$$

Insert it in (*):

$$\begin{aligned} & \frac{C}{N} \int dx dy \int_0^\infty \frac{dr}{r^5} \int dz \chi_{(r,z)}(x) \chi_{(r,z)}(y) \\ & \quad \times \langle \mathcal{U}_N(t)\Omega, a(\omega_{t,x})a(\omega_{t,y})a(1-\omega_{t,y})a(1-\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle \\ & = \frac{C}{N} \int_0^k dr \int dx dy dz \dots + \frac{C}{N} \int_k^\infty dr \int dx dy dz \dots \end{aligned}$$

Fefferman – de la Llave representation

$$(*) \quad \frac{1}{N} \int dx dy \frac{1}{|x-y|} \langle \mathcal{U}_N(t)\Omega, a(\omega_{t,x})a(\omega_{t,y})a(1-\omega_{t,y})a(1-\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle$$

Smooth version of **Fefferman–de la Llave representation** for the Coulomb potential:

$$\frac{1}{|x-y|} = C \int_0^\infty \frac{dr}{r^5} \int dz \chi_{(r,z)}(x) \chi_{(r,z)}(y), \quad \chi_{(r,z)}(x) = e^{-|x-z|^2/r^2}$$

Insert it in (*):

$$\begin{aligned} & \frac{C}{N} \int dx dy \int_0^\infty \frac{dr}{r^5} \int dz \chi_{(r,z)}(x) \chi_{(r,z)}(y) \\ & \quad \times \langle \mathcal{U}_N(t)\Omega, a(\omega_{t,x})a(\omega_{t,y})a(1-\omega_{t,y})a(1-\omega_{t,x})\mathcal{U}_N(t)\Omega \rangle \\ & = \frac{C}{N} \int_0^k dr \int dx dy dz \dots + \frac{C}{N} \int_k^\infty dr \int dx dy dz \dots \end{aligned}$$